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1 **An Investigation of Particle Breakage under One-Dimensional Compression of Sand Using**  
2 **X-ray Micro-Tomography**

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8 **Abstract.** Particle breakage alters particle-scale properties of granular soils including particle size,  
9 shape, and contact conditions, and changes macro-scale properties including soil compressibility,  
10 shear strength, and permeability. This study monitors the crushing of natural quartz sands under  
11 one-dimensional compression with in-situ X-ray tomography, i.e., X-ray scans during loading. We  
12 use the assembly-scale and particle-scale images to characterise particle failure patterns, e.g.,  
13 chipping, major splitting and comminution. Image processing and analysis enable us to determine  
14 the failure patterns around the yield stress, and the influence of initial density and particle  
15 morphology on the particle survival probability. We further quantify the degree of particle  
16 breakage with fractal dimension, breakage factor, and specific surface. Particle shape and  
17 coordination number both show a scale-dependent evolution pattern.

18  
19  
20 **Keywords:** fractals; laboratory tests; particle crushing/crushability; particle-scale behaviour;  
21 sands

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24 **1. Introduction**

25 Particle breakage in granular materials is relevant to many geotechnical infrastructures,  
26 such as offshore piles (Randolph 1988), high rock-fill dams (Sowers et al. 1965) and railway  
27 ballast (Indraratna et al. 1998). In a granular assembly, individual particles serve as the structural  
28 units with inherent properties, e.g., mineralogy, morphology, and strength. A particle will break  
29 into a few pieces once its internal stress reaches the particle failure strength. At the micro-scale,  
30 particle breakage changes inter-particle contact conditions, and material fabrics, leading to local  
31 stress redistribution. At the macro-scale, particle breakage alters soil compressibility (Pestana and  
32 Whittle 1995), shear strength (Marsal 1967), dilatancy (Yamamuro and Lade 1996) and  
33 permeability (Lade et al. 1996).

34 X-ray micro-computed tomography ( $\mu$ CT) is a non-destructive imaging tool, i.e., it does  
35 not alter the internal structure during the testing process. The high resolution of  $\mu$ CT images has  
36 allowed investigations into particle-scale properties of sands, including inter-particle contact  
37 conditions (Hasan and Alshibli 2010; Andò et al. 2013; Druckrey et al. 2016; Wiebicke et al. 2017;  
38 Afshar et al. 2018), particle morphology/material fabrics (Fonseca et al. 2012 and 2013), the  
39 evolution of particle breakage (Cil and Alshibli 2014; Alikarami et al. 2014; Karatza et al. 2017;  
40 Guida et al. 2018; Alam et al. 2018), and particle kinematics (Hall et al. 2010; Andò et al. 2012).

41 This experimental study uses X-ray  $\mu$ CT to characterise the particle breakage behaviour  
42 within an assembly of sand particles which is subjected to one-dimensional compression. We  
43 collect five scans at predetermined stress levels and apply image processing and analysis methods  
44 to quantify survival probability, failure pattern, and the evolution of particle size, shape, and  
45 coordination number.

46

## 47 2. Experimental Methods

### 48 2.1. Mini Oedometer Apparatus

49 The mini oedometer apparatus, which was fabricated to allow high-resolution laboratory  
50 CT scanning during the loading process, has a two-layer specimen chamber made of X-ray  
51 transparent carbon fibre reinforced polymer with high stiffness and strength, as shown in Fig. 1a.  
52 The external chamber bears the reaction force, while the internal floating ring provides lateral  
53 confinement and remains mobile relative to both loading platens. The loading platens are made of  
54 ceramic with low density to reduce the X-ray attenuation contrast and high hardness to avoid  
55 indentation of the platens. The cell body includes a stepping motor to apply the vertical load, and  
56 a load cell and an LVDT to monitor axial load and displacement, respectively. The controlling and  
57 monitoring system of the apparatus is located outside of the scanner and is not interrupted during  
58 scanning.

### 59 2.2. Materials and Tests

60 The Leighton Buzzard sand (LBS) with a particle diameter range of 0.6 – 1.18 mm was  
61 used in this study. LBS is mainly composed of quartz and contains relatively round and smooth  
62 particles. A detailed investigation of the single particle crushing behaviour of LBS particles was  
63 previously carried out by the authors (Zhao et al. 2015).

64 Loose and dense specimens of a diameter 6 mm were prepared within the floating ring  
65 through slow deposition and layer-by-layer compaction, respectively. The specimens were then  
66 compressed at a constant rate of 1.8 mm/h and scanned at five predetermined stress levels: S1 (1  
67 MPa) – at the initial loading, S2 (10 MPa) – around the yield stress, and S3-S5 (14, 19 and 28 MPa)  
68 – during post-yield compression, as shown in Fig. 1b. The yield stress  $\sigma_y$  is defined as the point

69 with the maximum curvature in the  $e$ -log  $\sigma$  curve for the loose specimen and is approximately  
70 taken as 10 MPa. For each scan, we paused the loading and waited for about ten minutes to allow  
71 the specimen to stabilise.

72 In this study, we used a commercial scanner Phoenix v|tome|x located at General Electric  
73 Company, Shanghai, China, and a custom-built scanner RX-Solutions at Laboratoire 3SR,  
74 Grenoble, France. They both have a high-power cone-beam source, a high-resolution detector, and  
75 a precision rotation stage. We performed tests on loose and dense specimens, LBS-L-S and LBS-  
76 D-S, in Shanghai with a voxel size of 4.0  $\mu\text{m}$ . A duplicate test is performed on a loose specimen,  
77 LBS-L-G, in Grenoble with a voxel size of 5.5  $\mu\text{m}$ .

### 78 *2.3. Image Processing and Analysis*

79 The processing and analysis of tomography images include three major steps. First, we  
80 classified the voxels of the filtered images into different phases, i.e., air, carbon fibre, sand particles,  
81 and ceramic platens, based on the threshold values determined by the multiple Otsu method (Liao  
82 et al. 2001). Second, we identified individual particles and fragments and labelled them with  
83 different integers. Finally, we quantified the 3D morphology of individual particles/fragments with  
84 a series of characteristic parameters. More details of the image processing can be found in Zhao  
85 and Wang (2016).

86 Figure 2 shows the labelled particles in LBS-L-S at S1 and S5. The image processing  
87 resulted in an overall satisfactory identification of individual particles before and after extensive  
88 fragmentation. However, difficulties arose in segmenting tiny fragments and accurately evaluating  
89 their shape parameters. Some tiny fragments remain very close to each other with a separation  
90 distance less than a few voxels especially near contacts, as shown by the grey-value vertical slices  
91 at high-stress levels in Fig. 3. Therefore, we excluded fragments with a volume smaller than 0.001

92 mm<sup>3</sup> (about 6000 voxels) from further analysis. Note that the total volume of particles has a relative  
93 error < 2.1% for all stress levels.

94 Following Zhao and Wang (2016), the particle size parameters include volume  $V$ , surface  
95 area SA and three principal dimensions ( $a \geq b \geq c$ ). Particle volume was used due to its lower  
96 dependency on particle shape. Two aspect ratios were used for a 3D particle, i.e., elongation index  
97 (EI=  $b/a$ ) and flatness index (FI=  $c/b$ ). The representative aspect ratio was defined as AR=  
98 (EI+FI)/2. The 3D mean-curvature roundness index  $R_M$  was evaluated at particle corners based on  
99 local mean curvatures. Particle compactness is quantified with two parameters, i.e., sphericity  $S =$   
100  $\sqrt[3]{36\pi V^2}/SA$  and convexity  $C_X = V/V_{CH}$ , where  $V_{CH}$  is the volume of particle's convex hull.

#### 101 2.4. Coordination Number

102 We measured the coordination number (CN) for each particle using the algorithm described  
103 in Andò (2013) and Andò et al. (2013). The algorithm's accuracy mainly depends on the accuracy  
104 of particle segmentation (Jaquet et al. 2013; Wiebicke et al. 2017). It should be mentioned that the  
105 algorithm considers multiple contact points between two particles (if existing) as one contact point,  
106 and the contacts with the loading platens and floating oedometer ring are not considered.

107

### 108 3. Experimental Results

#### 109 3.1. Failure Patterns at Assembly- and Particle-Scale

110 We describe particle-scale failures with chipping, splitting, and comminution (see  
111 illustrating sketches in Fig. 3). Chipping refers to minor parts split away from the particle, while

112 splitting indicates the particle splits into two or more major pieces. Comminution describes the  
113 extensive fragmentation resulting from a series of failure events.

114 Figure 3 shows the scan-by-scan comparison between the vertical slices of LBS-L-S at  
115 different stress levels, which identifies two chipping failures (S1-S2), two splitting failures of  
116 elongated particles and three chipping failures (S2-S3), and extensive fragmentation, and  
117 comminution (S3-S5). The fragments remained in proximity to each other, especially for the dense  
118 specimen.

119 Figure 4 illustrates the particle-scale failure patterns with four particles. Particle A  
120 experiences a progressive failure as the stress increases. First, a fracture initiates from the lower  
121 contact, propagates to a close concave feature and results in a chipping failure (A-S3). The  
122 conchoidal fracture occurs since the quartz particle has no cleavage. Then, a splitting failure occurs  
123 by a fracture connecting two contact points and generates a few fragments around the contact  
124 points (A-S4). Finally, the comminution of existing fragments generates many fragments with  
125 varying particle size and shape (A-S5).

126 Particles B and C with elongated morphology fail in bending as suggested by the cantilever  
127 curls. The fractures initiate at the tensile sides (marked by an arrow), then slow down and change  
128 direction at the tensile/compression boundary. A concave feature serves as a surface weakness and  
129 the fracture initiation point (C-S2). A small amount of LBS particles (<1%) have initial internal  
130 voids, which may also affect the fracture patterns (particle D). These typical examples demonstrate  
131 the varying particle failure patterns, which differ from the single particle crushing pattern under  
132 two flat platens (Zhao et al. 2015).

### 133 3.2. Volume-Change Ratio of Crushed Particles at S2

134 Visual inspection method (VIM) identified 68 initial particles that crushed into 177  
135 fragments at S2 from the labelled images of three specimens, as illustrated in the Appendix. The  
136 vertical stress values at S2 are 9.51 MPa, 9.70 MPa and 10.72 MPa for LBS-L-S, LBS-D-L, and  
137 LBS-L-G, respectively. Most crushing events (~ 85%) generate two or three fragments. A volume-  
138 change ratio is defined as  $R_v = (V_0 - V_1)/V_0$ , where  $V_0$  is the mother particle volume, and  $V_1$  is the  
139 volume of its largest fragment. Fig. 5a illustrates the volume-change ratio with the 3D view of four  
140 particle failure patterns. Fig. 5b shows the cumulative distribution of  $R_v$  for loose and dense  
141 specimens. The loose specimens have about 10% particles crushed with  $R_v$  up to 0.65, while the  
142 dense specimen has only about 4% particles crushed with  $R_v$  up to 0.23. Note that a fragment  
143 tracking algorithm is needed to analyse a large number of failure events in the post-yield scans.

### 144 3.3. Particle Survival Probability

145 We used the labelled images to inspect visually the intactness of initial particles in the  
146 following scans to determine their survival probability. This analysis considers only major splitting  
147 and excludes chipping failures. The reduction of survival probability can be described in two stages:  
148 Stage-1, a small percentage of particles crush up to 3% strain level (around  $\sigma_y$ ); Stage-2, the  
149 survival probability reduces linearly with the strain level, and about 50% particles fail at the highest  
150 strain level (Fig. 6a). Cil and Alshibli (2014) reported similar results in oedometer tests on brittle  
151 silica sand particles ( $d \sim 0.7$  mm) using X-ray CT images and discrete element simulation. Their  
152 numerical results agree well with our experiments for both loose and dense specimens (Fig. 6a).  
153 A similar level of particle failure percentage leads to a lower vertical strain in the denser specimen.  
154 The lower deformability in a denser specimen is mainly attributed to the fact that the particles

155 surrounding a crushing event have a higher chance to reconstitute a skeleton. The skeleton  
156 formation may also reduce the possibility of further comminution of the existing fragments (e.g.,  
157 A-S5 in Fig. 4).

158 We group the initial particles according to their shape parameters and evaluate their  
159 survival probability at the highest stress level (S5). Each group contains more than twenty particles  
160 for a reliable statistical analysis. The particles with lower aspect ratio, sphericity and convexity are  
161 more likely to experience major splitting failure (Fig. 6b). For example, the particles with AR=0.85  
162 and 0.6 have a survival probability of 69% and 22%, respectively. Conversely, the influence of  
163 particle roundness on failure probability is not obvious due to two major reasons: (1)  $R_M$  is an  
164 overall roundness index, which reflects poorly the morphology at contacts; and (2) chipping  
165 failures are not considered in evaluating survival probability.

#### 166 3.4. Particle Size Distribution

167 The cumulative particle size distribution (PSD) changes slightly from S1 to S4 and more  
168 extensively from S4 to S5 (Fig. 7a). The frequency PSDs from different scans intersect at the point  
169 of  $V=0.16 \text{ mm}^3$ , as shown in Fig. 7b, suggesting that the continuous fragmentation of various sized  
170 particles produces a balance of volume change at that point. The fragments smaller than the  
171 minimum initial particle ( $0.1 \text{ mm}^3$ ) form a fractal distribution:  $N(>V) \sim V^{-D/3}$ , where  $N(>V)$  is the  
172 number of fragments larger than  $V$  (Fig. 7c). The fractal dimension  $D$  increases almost linearly  
173 with the vertical stress in a semi-log plot after S2 (Fig. 7d).

#### 174 3.5. Breakage Factor and Specific Surface

175 The breakage factor  $B_f$ , which was defined by Nakata et al. (1999) as the volume ratio of  
176 particles smaller than the minimum particle size in the original specimen ( $0.1 \text{ mm}^3$ ), was adopted

177 to quantify the particle breakage degree, as shown in Fig. 8a. The curve shows a distinct bilinear  
178 pattern, with the turning point located around the yield stress. The dense specimen generates a  
179 smaller amount of fines and a lower  $B_f$  value at the maximum stress level than the loose specimens.

180 In addition to the breakage factor, the specific surface is another important parameter  
181 suitable for the measurement of particle breakage degree and often used in the constitutive  
182 modelling of crushable sands (e.g., McDowell et al. 1996; Russell, 2011). In this study, the specific  
183 surface is estimated based on  $\mu$ CT data as  $S_s = SA/(V \cdot \rho)$ , where  $\rho$  is the quartz density (2.65  
184 g/cm<sup>3</sup>). Similar to the approach adopted by Zhao et al. (2015), we used a fractal-based method to  
185 estimate the surface area of the fragments between 0.001 and 10<sup>-12</sup> mm<sup>3</sup> by assuming the sphericity  
186 to be the average sphericity of the fragments smaller than 0.1 mm<sup>3</sup> (Fig. 9c). These fragments  
187 contribute to about 10% of the total surface area at the maximum stress level. The three specimens  
188 have an initial specific surface  $S_{s0} \sim 3$  m<sup>2</sup>/kg, which is equivalent to a sphere with  $d = 0.75$  mm  
189 (Fig. 8b).  $S_s$  remains almost constant before S2 and increases gradually to about 1.5 $S_{s0}$  at the  
190 maximum stress level (about 2.8 $\sigma_y$ ).

### 191 3.6. Particle Shape

192 Particle shape shows a scale-dependent evolution: larger fragments evolve towards a cube  
193 with higher particle form and lower roundness, while smaller fragments are more flat (Fig. 9). The  
194 mean values of  $S$  and  $C_x$  remain almost constant for smaller fragments, due to the opposite trends  
195 of AR and  $R_M$ . This constant sphericity also exists for the fragments generated in single particle  
196 crushing tests (Zhao et al. 2015). We compare the 3D shape evolution with the 2D quantification  
197 results in a similar experiment on 1-2 mm LBS particles from Altuhafi and Coop (2011). Both 2D  
198 and 3D results reflect the reduction of particle form, sphericity, and convexity.

199 3.7. Coordination Number

200 The evolving particle size distribution alters the coordination number (CN) for different-  
201 sized particles. As shown in Fig. 10a, the average CN for particles of a given size increases  
202 significantly with the increasing particle breakage within the specimen, and such a trend is much  
203 more pronounced for larger particles. This is because a great number of smaller fragments  
204 generated fill the voids and surround the larger particles. At the highest stress level (i.e., S5), it is  
205 evident that the CN vs  $V$  plot reaches a fractal distribution:  $CN(>V) = 21V^{-0.33}$ . The fractal  
206 dimension equals 0.99, which is slightly larger than the fractal dimension for PSD ( $D = 0.9$ ). De  
207 Bono and McDowell (2016) reported a similar trend with a larger fractal dimension  $\sim 1.7$  from  
208 numerical simulations.

209 A large number of smaller fragments fill the voids and have low CN values. As a result,  
210 the average CN for all the particles throughout the compression process remains nearly constant  
211 at 6, as shown in Fig. 10b. To reflect better the influence of particle volume on the CN during the  
212 compression process, we plotted the volume-weighted average CN, which was calculated by  
213 multiplying each particle's CN by its volume (as a weighting factor), against the applied vertical  
214 stress in Fig. 10b. The curve shows again a bilinear pattern with a sharp increase of CN from S2.

215

216 **4. Conclusions**

217 This experimental study explores the particle breakage behaviour within an assembly of  
218 LBS particles under one-dimensional compression using X-ray  $\mu$ CT. The major thrust of this study  
219 lies in its comprehensive analysis of a variety of particle-scale properties based on  $\mu$ CT data,  
220 unveiling the microscopic origin of the macroscopic mechanical behaviour, particularly the

221 bilinear pattern of the stress-strain curve, yielding, and post-yield rapid compression. In detail, the  
222 following conclusions can be drawn from this study:

223 • In a sand assembly, particle morphology and multiple contacts lead to the complexity and  
224 variety of particle failure patterns. For example, elongated particles are more likely to fail with  
225 bending failures. Contacts and surface/internal flaws are prone to experience stress  
226 concentration and serve as crack initiation points.

227 • The consistent patterns of the fractal dimension, breakage factor, specific surface and volume-  
228 weighted CN versus the applied vertical stress demonstrate that these parameters work well to  
229 reflect the geometric nature of sand fragmentation and fabric evolution during particle  
230 breakage.

231 • A dense specimen has initially a lower failure probability and less extensive failure modes  
232 (smaller  $R_v$ ) than a loose specimen around the yield stress. The influence of initial density on  
233 particle survival probability diminishes after substantial particle breakage. However, a particle  
234 failure event leads to higher compression in a loose specimen and generates more fines through  
235 further comminution of existing fragments.

236 • Particle shape and coordination number both show a scale-dependent evolution. Larger  
237 particles evolve towards a cubic shape with high CN, while smaller fragments are flatter and  
238 floating with a small CN. The scale-dependent evolution reflects the different loading  
239 conditions and failure modes for coarse and fine particles.

240

241

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332 **Figure List**

- 333 • Fig. 1. The radiographs of specimen chamber (a) and the compression curves (b).
- 334 • Fig. 2. Labelled images for LBS-L-S at the initial loading state (S1) and the highest stress level  
335 (S5).
- 336 • Fig. 3. Grey-level vertical slices of LBS-L-S at multiple stress levels and of LBS-D-S at the  
337 highest stress level. Sketches illustrate different failure modes. Note: blue dots identify scan-  
338 by-scan failure events, and red circles indicate the particles used to illustrating failure modes.
- 339 • Fig. 4. Illustration of particle-scale failure modes. Particle A: progressive failure. Particles B  
340 and C: bending failures of elongated particles. Particle D: internal void related failure. Note:  
341 CL – cantilever curl; black circles mark fracture initiation locations.
- 342 • Fig. 5. (a) 3D view of failure patterns with varying volume-change ratios. (b) Cumulative  
343 failure percentage concerning  $R_v$  in the loose and dense specimens at S2 (~10 MPa).
- 344 • Fig. 6. (a) Evolution of survival probability with the increasing vertical strain. The results of  
345 the discrete element method DEM simulations on loose and dense specimens are from Cil and  
346 Alshibli (2014). (b) Influence of particle morphology on the survival probability at the highest  
347 stress level (S5).
- 348 • Fig. 7. Evolution of particle size distribution for LBS-L-S: (a) cumulative PSD; (b) frequency  
349 PSD; (c) fractal distribution of cumulative particle number. (d) Evolution of fractal dimension  
350 with the increasing vertical stress.
- 351 • Fig. 8. Evolution of the breakage factor (a) and the specific surface (b) with the increasing  
352 vertical stress. Note: the specific surfaces for a quartz sphere with varying diameters are  
353 marked by the blue lines in (b).

- 354 • Fig. 9. Scale-dependent evolution of particle shape parameters for LBS-L-S: (a) aspect ratio  
355 AR; (b) roundness  $R_M$ ; (c) sphericity  $S$ ; (d) convexity  $C_X$ . The data of G1 are 2D quantification  
356 results on LBS particles (1-2 mm) before and after a one-dimensional compression of 30 MPa  
357 (from Altuhafi and Coop, 2011). Note: arrows indicate the morphology evolution for initial-  
358 sized particles, and the dashed lines represent the trend curves for small fragments at the  
359 highest stress level.
- 360 • Fig. 10. Scale-dependent evolution of coordination number in LBS-L-S (a), and the evolution  
361 of average CN and volume-weighted average CN ( $V$ ) with the increasing vertical stress (b).  
362

1 **Appendix: Visual Inspection Method (VIM)**

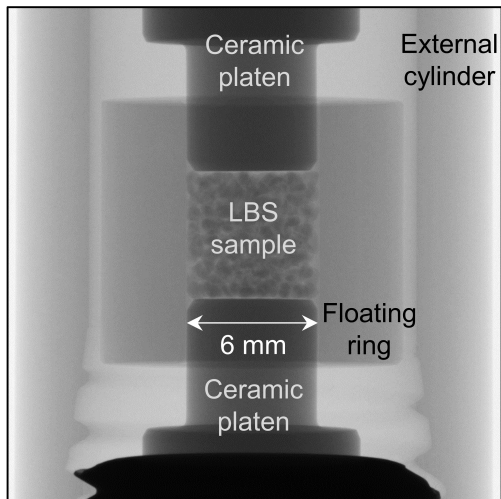
2 Visual inspection method (VIM) compares slice-by-slice the 3D labelled images in which particles  
3 are represented with distinctive integers to associate crushed particles with their fragments. For  
4 example, LBS-L-S initially has 308 particles in S1. VIM identified 28 crushed particles with 75  
5 associated fragments from S1 to S2. Fig. A-1 compares the 3D view of these particles before and  
6 after crushing. We quantified their failure patterns with the volume-change ratios ( $R_v$ ) using the  
7 3D volume of mother particles and the largest fragments. It shows that VIM performs well to  
8 identify crushing events in the current study. However, the efficiency of this method limits its  
9 application to a larger number of particles and fragments.

10

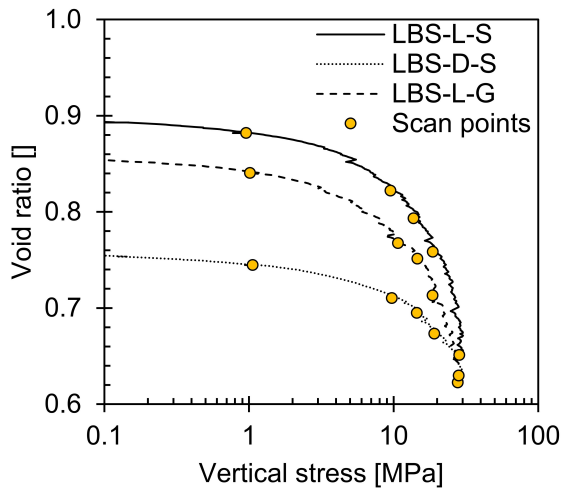
11

12

13 Fig. A-1. 3D view of particle crushing events from S1 to S2 in LBS-L-S together with volume-  
14 change ratios.



(a)



(b)

Fig. 1. The radiographs of specimen chamber (a) and the compression curves (b).

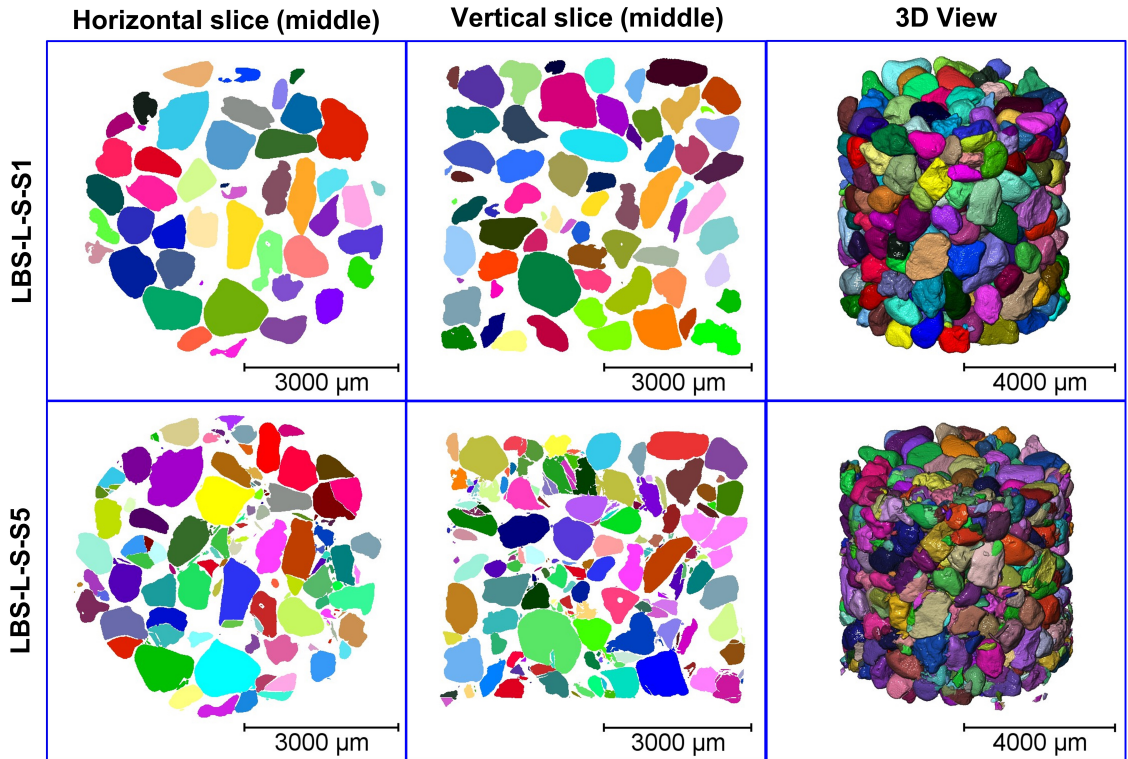


Fig. 2. Labelled images for LBS-L-S at the initial loading state (S1) and the highest stress level (S5).

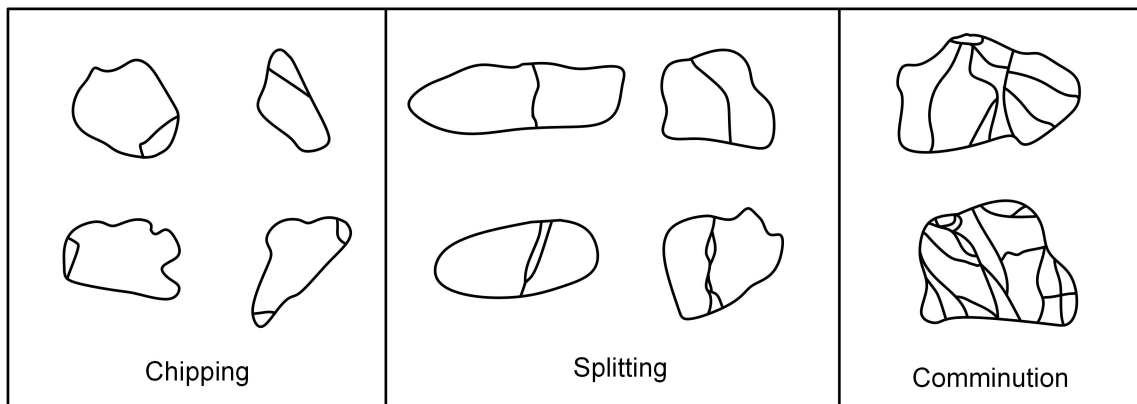
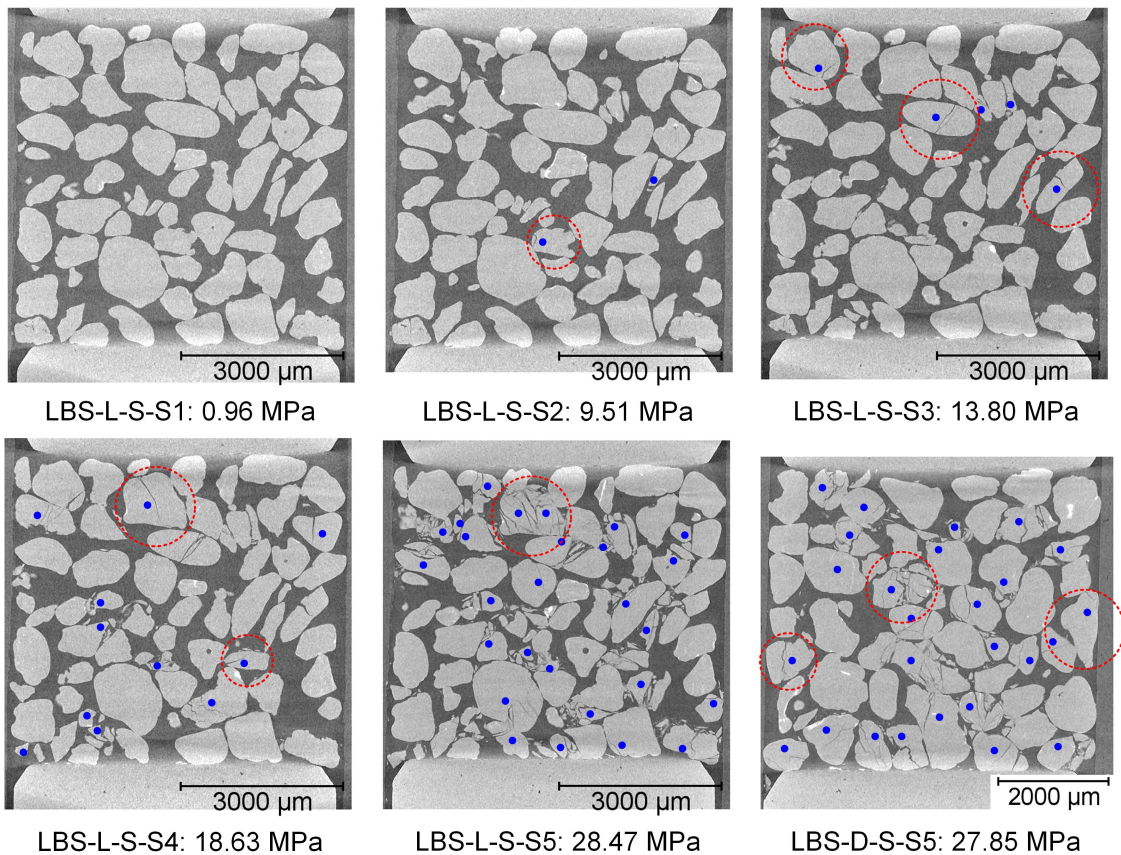


Fig. 3. Grey-level vertical slices of LBS-L-S at multiple stress levels and of LBS-D-S at the highest stress level. Sketches illustrate different failure modes. Note: blue dots identify scan-by-scan failure events, and red circles indicate the particles used to illustrating failure modes.

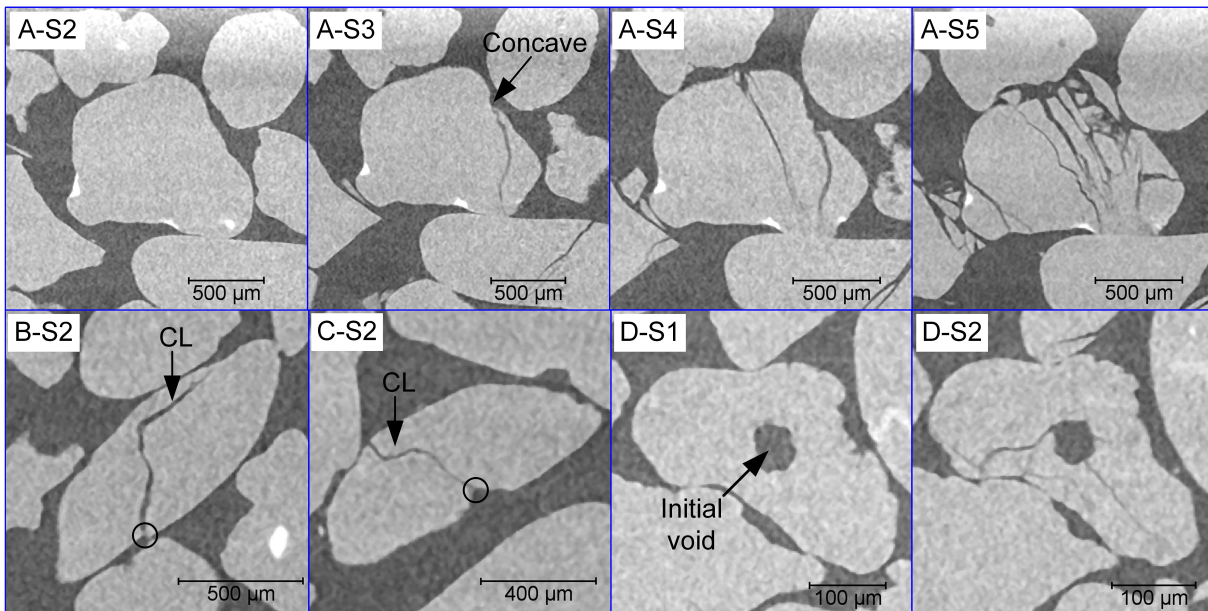
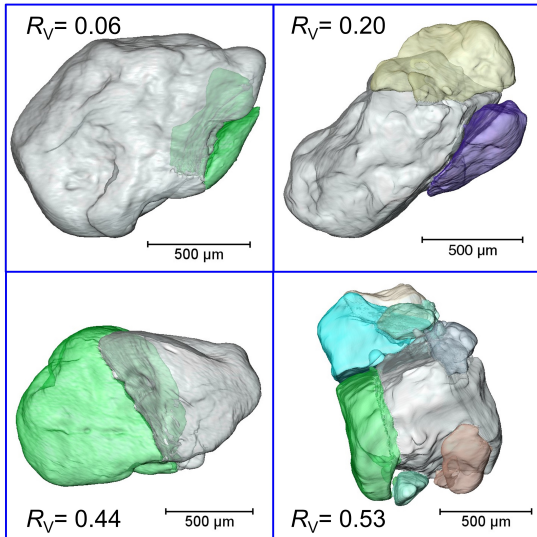
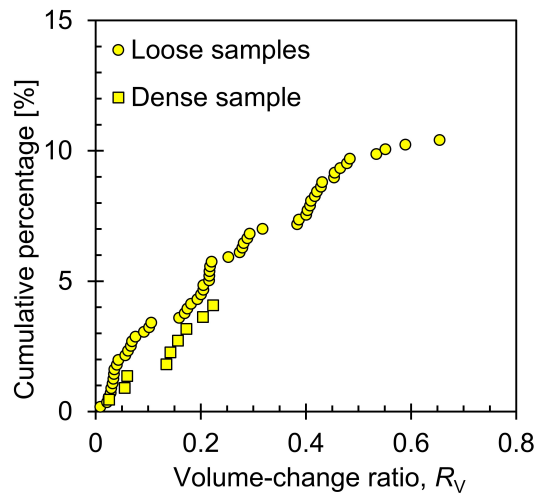


Fig. 4. Illustration of particle-scale failure modes. Particle A: progressive failure. Particles B and C: bending failures of elongated particles. Particle D: internal void related failure. Note: CL – cantilever curl; black circles mark fracture initiation locations.



(a)



(b)

Fig. 5. (a) 3D view of failure patterns with varying volume-change ratios. (b) Cumulative failure percentage concerning  $R_V$  in the loose and dense specimens at S2 ( $\sim 10$  MPa).

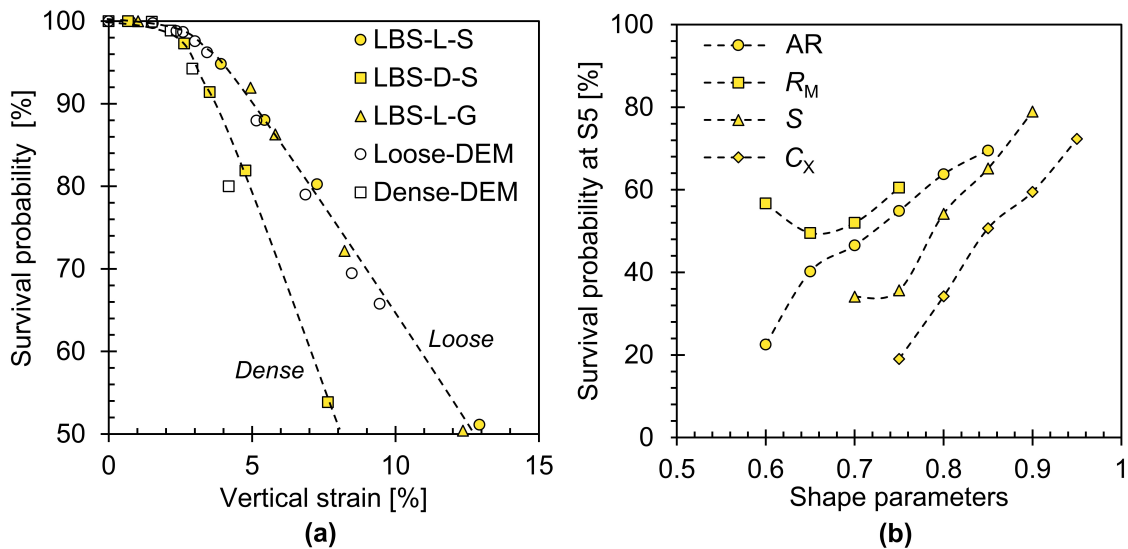
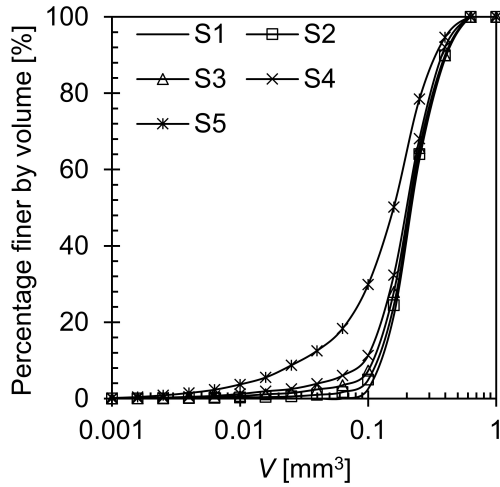
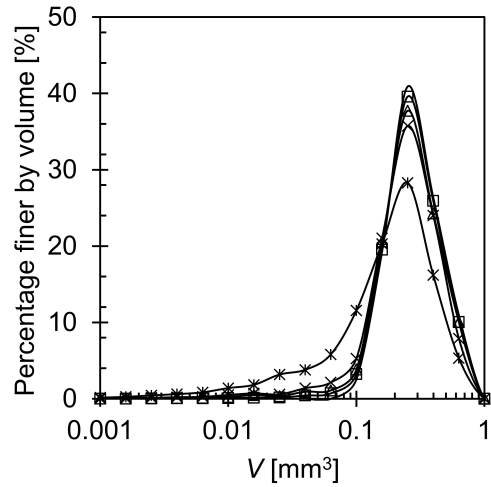


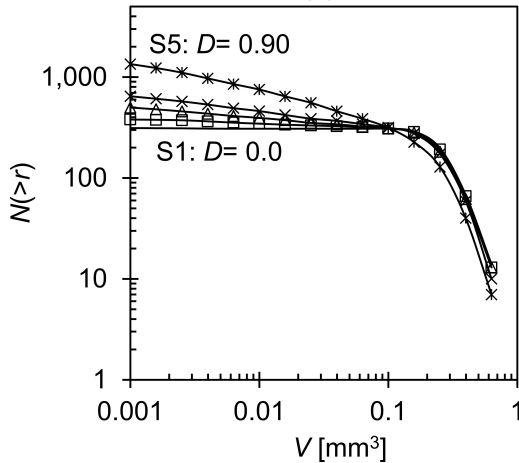
Fig. 6. (a) Evolution of survival probability with the increasing vertical strain. The results of the discrete element method DEM simulations on loose and dense specimens are from Cil & Alshibli (2014). (b) Influence of particle morphology on the survival probability at the highest stress level (S5).



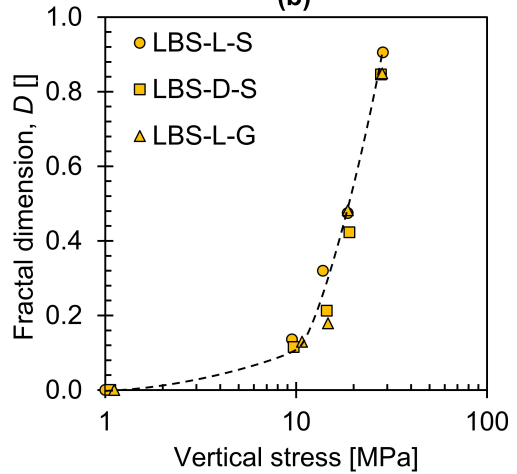
(a)



(b)



(c)



(d)

Fig. 7. Evolution of particle size distribution for LBS-L-S: (a) cumulative PSD; (b) frequency PSD; (c) fractal distribution of cumulative particle number. (d) Evolution of fractal dimension with the increasing vertical stress.

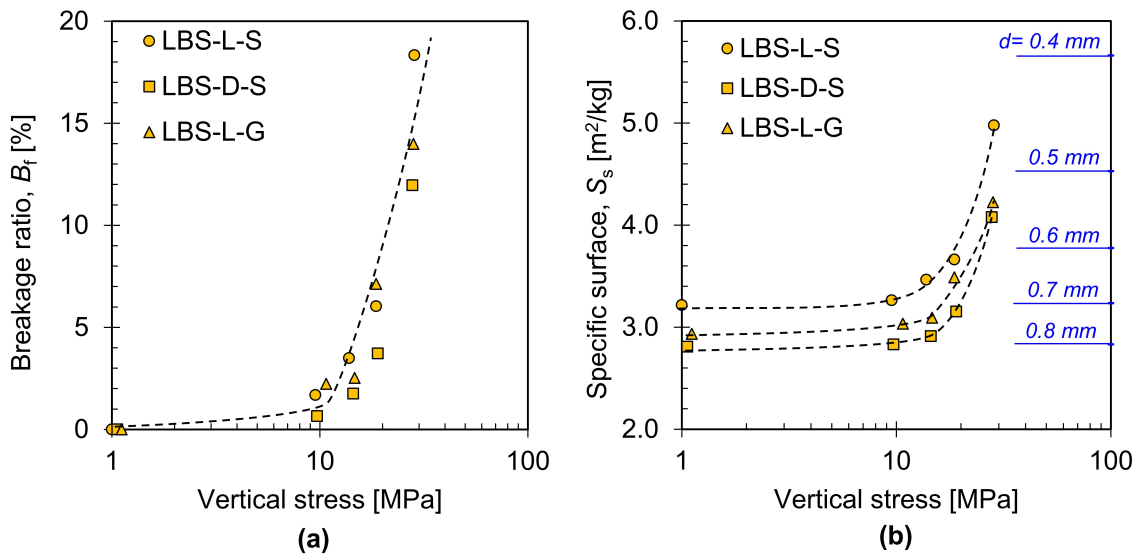


Fig. 8. Evolution of the breakage factor (a) and the specific surface (b) with the increasing vertical stress. Note: the specific surfaces for a quartz sphere with varying diameters are marked by the blue lines in (b).

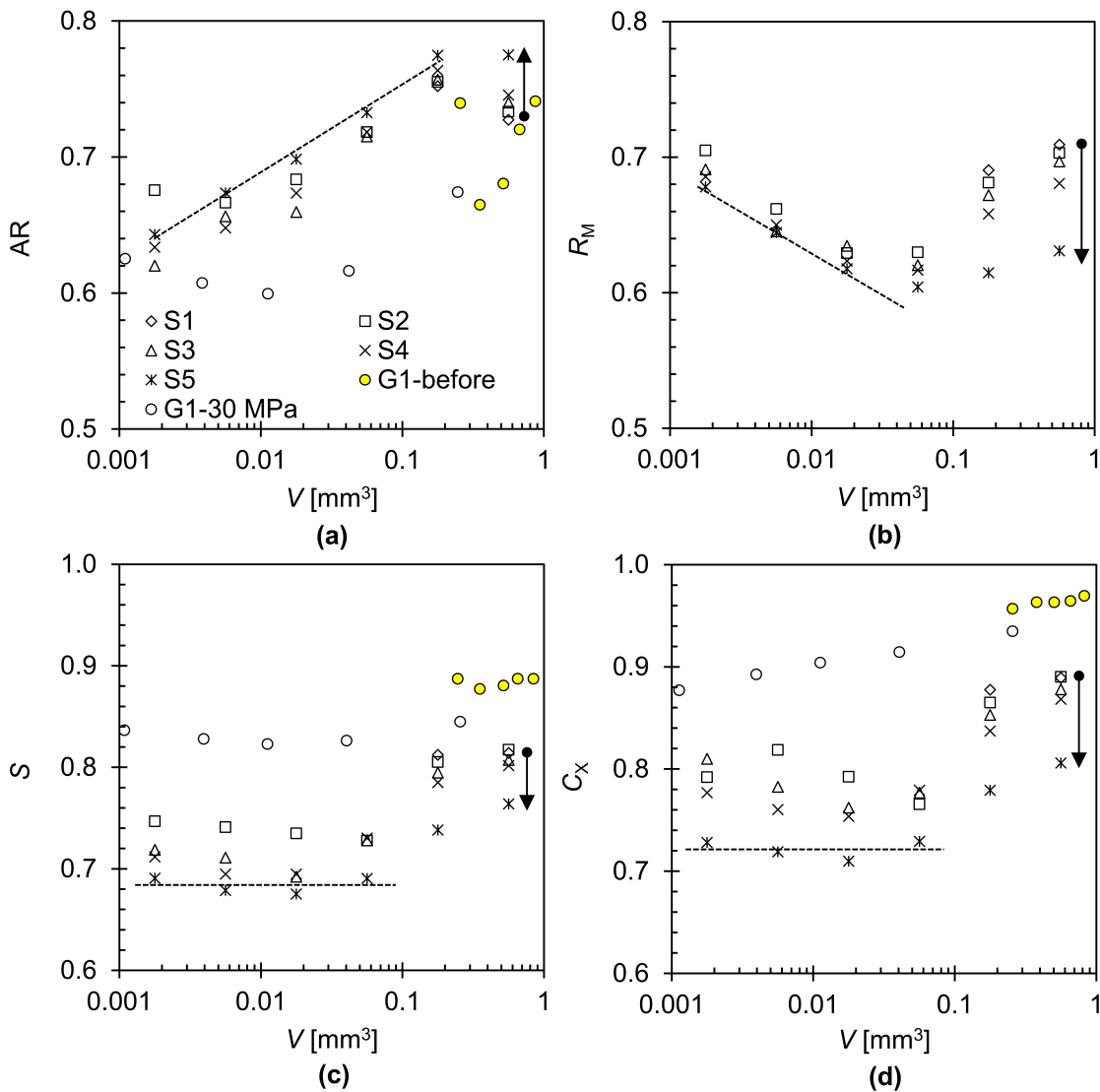


Fig. 9. Scale-dependent evolution of particle shape parameters for LBS-L-S: (a) aspect ratio AR; (b) roundness  $R_M$ ; (c) sphericity S; (d) convexity  $C_X$ . The data of G1 are 2D quantification results on LBS particles (1-2 mm) before and after a one-dimensional compression of 30 MPa (from Altuhafi & Coop, 2011). Note: arrows indicate the morphology evolution for initial-sized particles, and the dashed lines represent the trend curves for small fragments at the highest stress level.

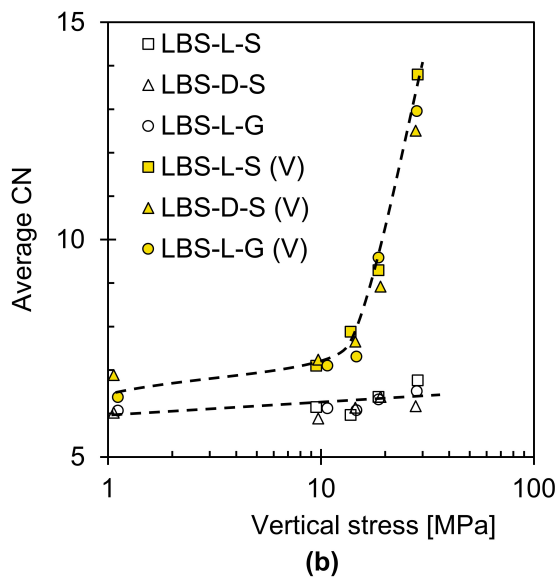
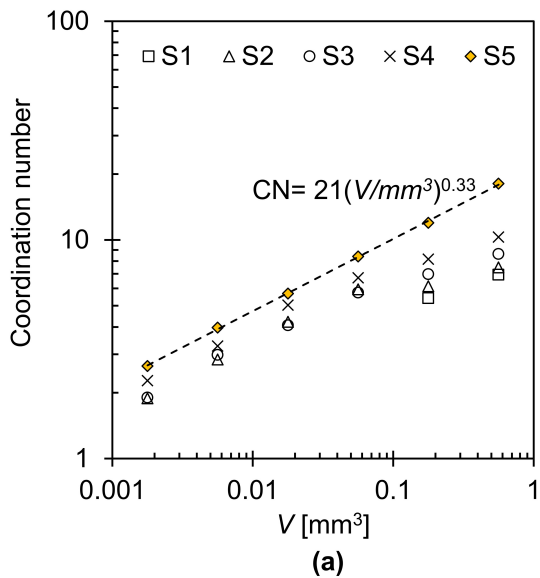


Fig. 10. Scale-dependent evolution of coordination number in LBS-L-S (a), and the evolution of average CN and volume-weighted average CN (V) with the increasing vertical stress (b).

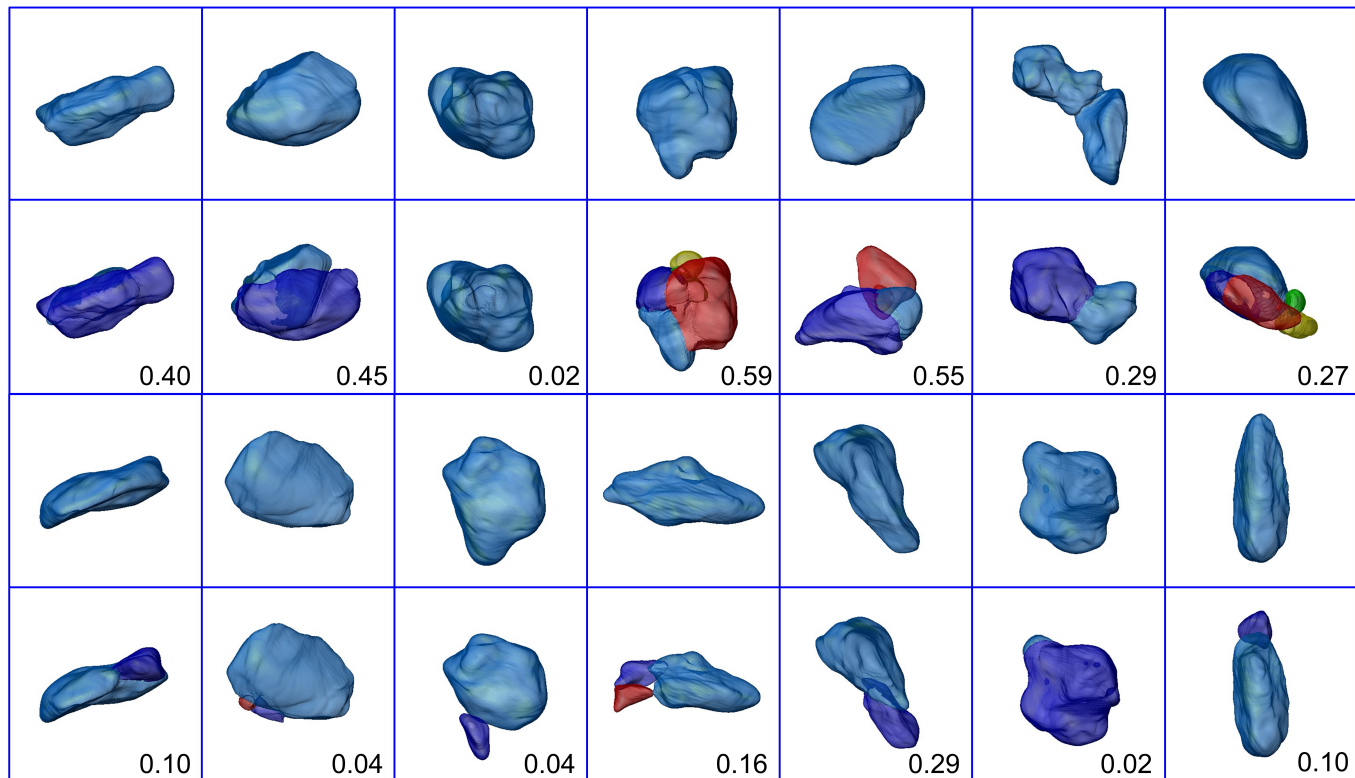


Fig. A-1. 3D view of particle crushing events from S1 to S2 in LBS-L-S together with volume-change ratios.

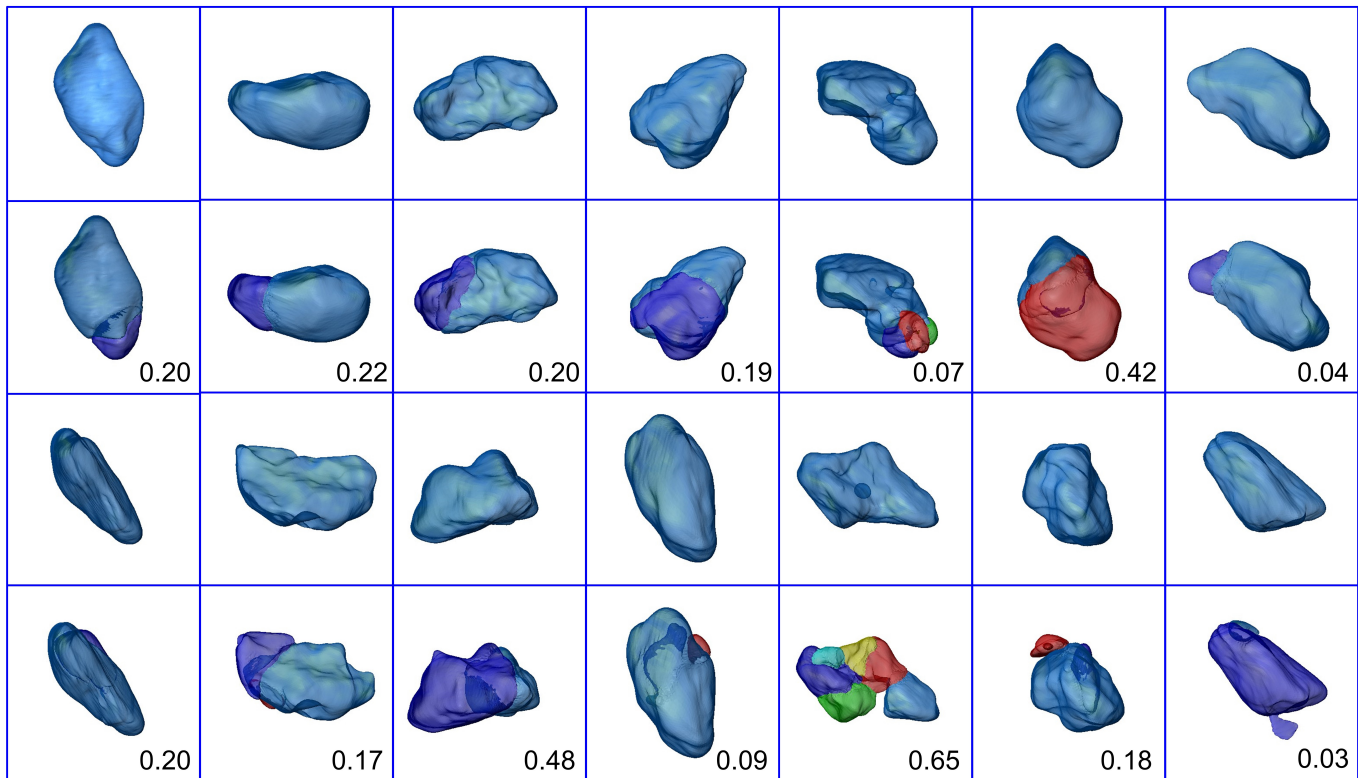


Fig. A-1. Continued.