Binary blazed reflection gratings


A reflection grating with a binary surface profile is presented that has high diffraction efficiency. The measured intensity for the +1st diffracted order was 77%. The binary grating is composed of a minilattice with feature sizes comparable with the wavelength of the incident light. The overall structure is designed in such a way that it imitates a conventional blazed grating. The grating also has interesting polarization properties. The main part of the TE-polarized light is diffracted into the 1st diffracted order, and most of the TM-polarized light remains in the 0th diffracted order. The measurements of the grating are compared with rigorous diffraction theory and found to be in reasonable agreement.

1. Introduction
Recently\textsuperscript{1-3} the manufacture of dielectric gratings has been proposed in which each period is composed of a minilattice with a variable duty cycle and binary surface profile. Such elements have an advantage compared with conventional blazed gratings, which are known to show such high diffraction efficiency that only one microlithographic fabrication step is necessary to produce such gratings. The minilattice is based on a zero-order grating, and its behavior is similar to a dielectric material. By varying the duty cycle of the minilattice, the effective refractive index of the microstructure can be changed. In this way a binary grating can be made that behaves like an artificial dielectric material with a distributed index. In fact similar ideas have arisen previously in the area of micrometer waves\textsuperscript{4-6}.

In this paper this idea is reapplied to create metallic binary blazed gratings as shown in Fig. 1 in principle. To understand the operation of these gratings, first we investigate the physics of the minilattices, which in this case are metallic zero-order gratings. In Section 2 the theory of such binary blazed gratings is treated briefly. The subject of Section 3 is the fabrication method of such a grating. In Section 4 the results of measurements are compared with rigorous diffraction theory. The disagreements between theory and experiment are discussed, and some possible applications of the grating are presented.

2. Theory
The theory of these devices was developed in a previous publication\textsuperscript{7}. Therefore in this section only a cursory version of the approximate model is presented so that it is possible to understand the grating operation.

In Fig. 1 a single period of such a binary blazed grating is shown. The height of the grating structure is \( h \), and the overall grating period is \( d_1 \), which corresponds to a deflection angle \( \sin(\phi_1) = d_1/\lambda \) for the first order. The grating is composed of a minilattice with \( M \) binary (rectangular) grooves and a period \( d_2 \). Each of these grooves has a variable width \( d_{3,m} \), where \( m \) is the index number of each individual microstructure; \( d_2 \) is assumed to be less than the wavelength of the incident light, so that all nonzero diffraction orders of the microstructure are evanescent. The refractive index above the grating is assumed to be \( n_1 = 1 \). The grating is designed to imitate a highly efficient blazed grating. To understand the operation of such a complicated structure, one must first understand the operation of metallic zero order gratings (Fig. 2), i.e., a minilattice with constant groove width \( d_3 \). If a plane wave is normally incident on such a grating, besides absorption, all the energy of the input beam is reflected into the zeroth diffraction order. The duty cycle \( t \) of this
Grating is defined as

\[ t = \frac{d_2 - d_3}{d_2}, \]  

which is the ratio of the filled metallic part of the microstructure to one period. The basic idea of this paper is the following: The phase of the reflected wave depends on the duty cycle of the minilattice period, and therefore by varying the duty cycle of the minilattice, one can fabricate a macroscopic blazed grating or more generally a kinoform. For a grating period \( d_2 = 9 \ \mu\text{m} \) and the height is \( h = 5.3 \ \mu\text{m} \), the refractive index of the grating is \( n = 0.22 + i6.71 \), the wavelength is \( \lambda = 10.6 \ \mu\text{m} \). The wave is TE polarized and normally incident.

3. Manufacturing the Grating

For the fabrication of the grating a minimum feature size of 1 \( \mu\text{m} \) was assumed, which is limited by the laser-pattern generator available to us. For fabrication a fused-silica mask blank, with a 0.1-\( \mu\text{m} \)-thick chromium layer, was coated with a 0.5-\( \mu\text{m} \)-thick layer of photoresist and was exposed with the laser-pattern generator (LPG-15P from Micronics) with a He–Cd laser. After resist development and chromium etching the fused silica was dry etched with reactive ion-etching equipment, PLASMA-LAB from Plasma Technology, with a CHF\(_3\) process at 25-mTorr pressure and 195-W high-frequency power. Because of insufficient etch selectivity between the SiO\(_2\) substrate and the photoresist layer, the maximum grating depth achievable was found to be \( \sim 2 \ \mu\text{m} \).

For the desired gratings, depths of \( \sim 5 \ \mu\text{m} \) are required, and therefore the depth was increased by a self-alignment technique. Using this technique, one coats the grating again with photoresist and exposes it from the back side. Developing and reactive ion etching are repeated until the desired depth is reached.

The last step in the process is cleaning the grating by removing the chromium followed by metallic coating (aluminum) by evaporation under several oblique angles, so that the perpendicular walls of the grating are also well coated with metal. Unfortunately the layer thickness achievable by this technique was not completely sufficient, which was observed by electrical conductivity measurements. The conductivity parallel and perpendicular to the grating lines shows a ratio of 50:1, which indicates a very thin coating on the perpendicular walls. Therefore some penetration of the incident light into the substrate might
The achieved depth of the grating was \( \approx 5.3 \, \mu m \) with an accuracy of 0.3 \( \mu m \) and was measured with an accuracy of 0.3 \( \mu m \) with a Linnik microscope. The grating period was 90 \( \mu m \). The number of minilattices was \( M = 10 \). In Table 1 the parameters of the grating are given. The feature sizes were estimated by a scanning electron microscope. Because of underetching feature sizes appear that are less than the rated minimum feature of 1 \( \mu m \). The first column denotes the particular \( m \)th microstructure being discussed. Columns 2, 3, and 4 are for the manufactured grating. Columns 5, 6, and 7 are for the proposed calculated grating with a 1-\( \mu m \) feature size. In column 2 left indicates the left starting point of a feature, and in column 3 right indicates the right end point of a feature. Column 4 lists the duty cycle in each minilattice. In Fig. 4 a scanning electron micrograph of the manufactured grating is shown.

### Table 1. Shape of the Minilattices of the Grating

<table>
<thead>
<tr>
<th>Minilattice</th>
<th>Manufactured Grating</th>
<th>Calculated Grating</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Left (( \mu m ))</td>
<td>Right (( \mu m ))</td>
</tr>
<tr>
<td>1</td>
<td>0</td>
<td>0</td>
</tr>
<tr>
<td>2</td>
<td>9</td>
<td>9.7</td>
</tr>
<tr>
<td>3</td>
<td>18</td>
<td>18.7</td>
</tr>
<tr>
<td>4</td>
<td>27</td>
<td>27.7</td>
</tr>
<tr>
<td>5</td>
<td>36</td>
<td>37.8</td>
</tr>
<tr>
<td>6</td>
<td>45</td>
<td>47.8</td>
</tr>
<tr>
<td>7</td>
<td>54</td>
<td>57.3</td>
</tr>
<tr>
<td>8</td>
<td>63</td>
<td>66.8</td>
</tr>
<tr>
<td>9</td>
<td>72</td>
<td>76.3</td>
</tr>
<tr>
<td>10</td>
<td>81</td>
<td>85.8</td>
</tr>
</tbody>
</table>

*Columns 2–4 refer to the experimentally realized component, and columns 5–7 describe the design. Apparently during the microlithographic fabrication some underetching systematically reduced all duty cycles.

In column 2 left indicates the left starting point of a feature, and in column 3 right indicates the right end point of a feature. Column 4 lists the duty cycle in each minilattice. In Fig. 4 a scanning electron micrograph of the manufactured grating is shown.

### 4. Measurements and Comparison with Theory

The diffraction efficiencies of the gratings were measured with a CO\(_2\) laser at a 10.6-\( \mu m \) wavelength. The incident intensity of the CO\(_2\) laser was \( \approx 4 \, W/cm^2 \). To obtain polarized light, a conventional Brewster arrangement was used.

The refractive index of aluminum at \( \lambda = 10.6 \, \mu m \) is \( n = 26.6 + i96.6 \),\(^1\) which is the refractive index of a nearly perfect conductor. For a comparison of the measurements two different rigorous diffraction theories were used:

(a) The differential method:\(^8\): Because of the numerical problems in our implementation of the differential method it was not possible to use the refractive index above. Instead the arbitrary index \( n = 0.22 + i6.7 \) was used, which is also the refractive index of a good metal. For the calculations 61 diffracted orders were included.

(b) The modal method:\(^12\)-\(^14\): The grating was assumed to be perfectly conducting. The electric/magnetic field inside the rectangular grooves can be expressed with modes. In the calculations four modes were included in every groove. In the calculation 61 diffracted orders were taken into account.

In Fig. 5 diffraction efficiencies of the +1st, +10th, −9th, −10th diffraction order are shown as a function of the angle of incidence. The solid (dotted) curves are calculated with the differential method (modal method). Besides some differences for the −9th and −10th diffraction orders at high angles of incidence the agreement is good. The maximum of the diffraction efficiency of the +1st order is 84% for the differential method and 88% for the modal method.

The metallic coating of the considered grating is not as good a conductor as a perfect metal but a better conductor than the grating calculated with the differential method. Therefore we assume that the diffraction efficiencies can be calculated by both methods. The errors from the incorrect index are relatively small.

In Fig. 6 the diffraction efficiencies for normal incidence and TE-polarized light are shown (a) with a linear and (b) with a logarithmic scale. These theoretical estimations were produced with the difference...
Fig. 6. Measured and calculated diffraction efficiencies for normal incidence and TE polarization. The grating period is $d_1 = 90 \mu m$, the wavelength is $\lambda = 10.9 \mu m$, and the height is $h = 5.3 \mu m$. The period of the $M = 10$ minilattice is $d_2 = 9 \mu m$. The minilattice parameters are described in detail in Table 1. The efficiencies are shown (a) with linear and (b) with logarithmic scaling.

Fig. 7. Diffraction efficiency of the +1st diffracted order for normal incidence for the grating in Fig. 4. But now the polarization angle $\omega$ is changed from 0 (TE polarization) to 90 deg (TM polarization).

The incident beam relative to the grating vector $K$ is presented. The calculated diffraction efficiencies of the +1st diffracted order are 88% (TE polarization) and 0% (TM polarization). The diffraction efficiencies of the 0th order are 1% (TE polarization) and nearly 100% (TM polarization). These diffraction efficiencies were calculated by the modal method. It is expected that the diffraction efficiency $\eta$ depends sinusoidally on the polarization angle $\omega$ (Ref. 15):

$$\eta(\omega) = \eta_{TE} \cos^2(\omega) + \eta_{TM} \sin^2(\omega),$$

where $\eta_{TE}$ ($\eta_{TM}$) indicates the efficiency of a particular diffraction order at normal incidence for TE (TM) polarization. The curves belong to the theoretical predictions. The great differences between theoretical and measured values are striking.

The measured efficiency $\eta_{TM}$ of the +1st order is ~2%, whereas $\eta_{TM}$ of the 0th order is ~50%. (The theoretical prediction is nearly 100%). From Fig. 7 we can see that the grating can be used as a polarizer, which diffracts most of the light for TM polarization into the 0th diffracted order and most of the light for TE polarization into the +1st diffracted order. When the polarization angle $\omega$ is changed, variable amounts of energy can be put in the 0th and +1st diffracted orders.

Now a set of measurements is described where the angle of incidence is varied (Fig. 8). Variation of the angle of incidence in the $xOz$ plane is called classical diffraction, and variation of the angle of incidence in the $zOy$ plane is called conical diffraction.

At first the case of classical diffraction is considered. The diffraction efficiencies as a function of the angle of incidence are examined. As has been said the grating period $d_1 = 90 \mu m$ and the number of minilattices $M = 10$. The incident wavelength $\lambda = 10.6 \mu m$. The variation of the incident beam is in the $xOz$ plane. In Fig. 8 the diffraction efficiencies of the +9th, +1st, 0th, -9th, and -10th diffraction orders are shown as functions of the angle of incidence. Measurements are compared with theoretical predictions calculated by the differential method. Under normal incidence 17 diffracted orders propagate. This fact can be calculated with the help of the...
Fig. 8. Difference between classical ($xOz$-plane) and conical ($zOy$-plane) variation of the angle of incidence.

grating equation:

$$(2\pi/\lambda)\sin(\varphi_L) = (2\pi/\lambda)\sin(\varphi) + (2\pi/d_1)L,$$  \hspace{1cm} (4)

where $\varphi$ is the angle of incidence and $\varphi_L$ is the angle of the $L$th diffracted order. At normal incidence, $\varphi = 0$, the $+10$, $+9$, $-9$, and $-10$th diffraction orders are evanescent. The $\pm 9$th diffracted order appears when $|\varphi| = 3.45$ deg, and the $\pm 10$th diffracted order occurs at $\pm 10.25$ deg. Now we return to the idea of the minilattices discussed in Section 2. The $\pm 1$st diffraction order of the minilattice appears when $|\varphi| = 10.25$ deg. The $\pm 10$th diffracted order of the superlattice and the $\pm 1$st diffraction order of the minilattice have the same diffraction angle and so correspond to each other. The duty cycle of the minilattice does not vary rapidly, and the diffraction spectrum roughly follows the spectrum for a perfectly periodic array. Therefore the strong increase of the $\pm 10$th diffracted orders can be explained by the fact that diffraction at the minilattice level occurs. The strong $-9$th diffracted order can be explained by the scatter of the strong $+1$st diffracted order by the minilattice.

By choosing a suitable minilattice period $d_2$, we can make the diffraction efficiency variation of the $+1$st diffracted order broader or narrower. As can be seen from Fig. 7 this function is almost flat until the strong diffracted orders start to propagate. At these angles the diffraction efficiency of the $+1$st diffracted order decreases very rapidly. Increasing the microstructure period $d_2$ decreases the width of the $+1$st diffracted order, and decreasing the size $d_2$ results in increasing the width of the function. Thus by changing the size $d_2$ we can design the angular selectivity of the $+1$st diffracted order. The diffraction efficiencies measured are always, especially for the $+1$st diffracted order, less than the theoretical predictions, probably because of losses in the grating. At an angle of $\approx 2$ deg the diffraction efficiency of the $+1$st diffracted order reaches 77%.

In Fig. 10 the diffraction efficiencies are shown as functions of the angle of incidence, but now for TM polarization. The variation of the incident beam is again in the $xOz$ plane. The $0$th diffracted order is the dominant order. The $+1$st diffracted order has almost vanished.

Now we describe a measurement in which the angle of the incident beam is varied in the $zOy$ plane as shown in Fig. 8, which is the special case of conical diffraction. For the conical angle of incidence of 0 deg the case corresponds to the case of classical diffraction. For the conical angle of incidence of 0 deg the $E$-field vector $E_{TE}$ is always parallel to the $x$ axis and the $E_{TM}$ vector is always in the $zOy$ plane. The $E_{TE}$ and $E_{TM}$ vectors are perpendicular to each other and to the $k$ vector, and they define a plane shown as a disk in Fig. 6. An approximate model exists for the general case of conical diffraction by a perfect lossless metal and is described in Ref. 16.

For the particular case examined here this model predicts that the diffraction efficiencies of all diffracted orders should vary as if the input light was normally incident but had a wavelength of $\lambda/\cos(\theta)$. This approximate model has been used for the calcula-

![Fig. 9](image_url)  
**Fig. 9.** Diffraction efficiencies as functions of the angle of incidence ($xOz$ plane) for TE polarization (classical diffraction). The symbols are measured values; the curves are calculated by the differential method.

![Fig. 10](image_url)  
**Fig. 10.** Diffraction efficiencies for TM polarization as a function of the angle of incidence ($xOz$ plane, conical diffraction). The symbols are measured values. A logarithmic scale is used.
5. Summary
We have presented a blazed reflection grating with a binary surface profile. The measured diffraction efficiency for TE polarization was 77%, whereas theory predicts an efficiency as great as 88% for the design used here. The grating can also be used as a polarizing beam splitter.

The measurements with the CO₂ laser were performed at the Lehrstuhl für Fertigungstechnik (by M. Geiger), University of Erlangen, Germany. The authors are grateful for the help and for the permission to use equipment from the Laser Research Cooperative, Erlangen, Germany. One of the authors (J. T. Sheridan) is currently funded by the European Community. Part of the research was also funded by the Deutsche Forschungsgemeinschaft under Sondriforschungsbereich 182. Cooperation with Micronic Laser Systems AB, Sweden regarding the Laser Pattern Generator LPG-15P is gratefully acknowledged.

References