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<tr>
<th><strong>Title</strong></th>
<th>Spread-space spread-spectrum technique for secure multiplexing</th>
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<tr>
<td><strong>Authors(s)</strong></td>
<td>Hennelly, Bryan M.; Naughton, Thomas J.; McDonald, John B.; Sheridan, John T.; Gopinathan, Unnikrishnan; Kelly, Damien P.; Javidi, Bahram</td>
</tr>
<tr>
<td><strong>Publication date</strong></td>
<td>2007-05-01</td>
</tr>
<tr>
<td><strong>Publication information</strong></td>
<td>Optics Letters, 32 (9): 1060-1062</td>
</tr>
<tr>
<td><strong>Publisher</strong></td>
<td>Optical Society of America</td>
</tr>
<tr>
<td><strong>Link to online version</strong></td>
<td><a href="http://dx.doi.org/10.1364/OL.32.001060">http://dx.doi.org/10.1364/OL.32.001060</a></td>
</tr>
<tr>
<td><strong>Item record/more information</strong></td>
<td><a href="http://hdl.handle.net/10197/3420">http://hdl.handle.net/10197/3420</a></td>
</tr>
<tr>
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<tr>
<td><strong>Publisher's version (DOI)</strong></td>
<td>10.1364/OL.32.001060</td>
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A novel technique for multiplexing complex images is proposed in which each image may be demultiplexed only if a set of random encryption keys is known. The technique utilizes the ability of the double random phase encoding method to spread a signal’s energy in both the space and the spatial frequency domains in a controlled manner. To multiplex, images are independently encrypted with different phase keys and then superimposed by recording sequentially on the same material. Each image is extracted by using the particular key associated with it. During decryption the energy from the other images is further spread, making it possible to minimize its effects by using suitable filters. Wigner analysis is applied to the technique, and numerical results are presented.

The standard DRPE method is well known. An input image wave field passes through a random diffuser (D1), an optical Fourier transform (OFT), a second diffuser (D2), and finally a second OPT. If the diffusers may be described as two statistically independent random phase functions, then the encrypted image may be described as a white process. Decryption involves passing the encrypted wave field back through the encryption system in which D2 is replaced by its conjugate D2* and D1 is simply removed. The effect of D2 is canceled, and D1* is not employed, since we are interested only in the intensity of the decrypted signal. Our multiplexing scheme comprises a DRPE and a recording material. A number of images are encoded and recorded sequentially by using the same reference beam but different diffusers. The effects of the diffusers on the spatial extent of an input signal and its spatial frequency extent as it evolves through these systems are well known. The WDF allows us to elegantly explain these effects for the purposes of this paper. The WDF of a 1D signal $u(x)$ is given as

$$\psi[u(x)](x,k_x) = \int u(x - \frac{\epsilon}{2})^* u^*(x + \frac{\epsilon}{2}) e^{-i2\pi k_x \epsilon} d\epsilon, \quad (1)$$

where $x$ denotes the spatial coordinate and $k_x$ denotes spatial frequency. In the case of a 2D signal the
WDF is four dimensional. Much of our analysis is limited to the 1D case, the extension to 2D being trivial. The WDF has many properties, only a few of which are critical in the present context. The first property is that of localization; in many practical problems it we assume that a signal is bounded within some finite region in the spatial and spatial frequency domains. The spatial extent, \( W \), and frequency extent, \( B \), are the minimum values such that

\[
u(x) = 0, \quad |x| > W, \quad \int \nu(x)e^{-i2\pi k_x x}dx = 0, \quad |k_x| > B.
\]

(2)

The second property relates to Parseval’s theorem; the total energy of the signal, \( E \), is

\[
E = \iiint \psi(u(x))\psi^*(x, k_x)dxdk_x.
\]

(3)

Applying Eqs. (2) and (3), the average energy density of the signal in phase space is given by \( E/WB \). The third property of interest is the product theorem of the WDF. When two signals are multiplied in space, their WDFs are convolved along the \( k_x \) axis:

\[
\psi(u(x)v(x))(x, k_x) = \int \psi(u(x))(x, k_x - k')\psi(v(x))(x, k_x)dk'.
\]

(4)

From Eqs. (2) and (4), the signal \( u(x)h(x) \) will have a spatial width equal to the overlapping spatial distributions of the individual signals. Furthermore, the resulting spatial frequency distribution will have a bandwidth equal to the sum of the bandwidths of the product signals, \( B_n + B_p \). The final property of interest is the relationship of the WDF to the Fourier transform (FT). If a FT is applied to a signal, the corresponding WDF rotates by \( \pi/2 \) rad: \( \psi(u(x))(x, k_x) \rightarrow \psi(u(x))(-k_x, x) \). In the case of an OPT with focal length \( f \) and wavelength \( \lambda \), we have the following relationship: \( \psi(u(x))(x, k_x) \rightarrow \psi(u(x))(-\sqrt{\lambda f_k}, x/\sqrt{\lambda f}) \). If Eq. (2) holds, then the output OPT bandwidth is \( B/\sqrt{\lambda f} \); i.e., the OPT does not simply rotate the WDF but also stretches it.

Turning to the DRPE system, and with the aid of Fig. 1 and Table 1, we describe the evolution of a signal’s spatial and spectral frequency distributions as it propagates through the system. The input signal has the Wigner volume chart shown in Fig. 1(i), where the signal is depicted having a width and bandwidth as defined in Table 1. The average energy density can be calculated at any plane as \( E/(\text{Width } x) \) (Width \( k_x \)). If our image signal has been coded for efficient storage, then it will seem random in nature, and we may assume that it fully occupies the volume shown in Fig. 1(i). D1 and D2 are assumed to be lossless and are designed to have the well-specified widths and bandwidths listed in Table 1. When the signal passes through D1, the product signal has an increased bandwidth, but the total energy must remain the same. Therefore the average energy density is reduced by a factor proportional to the spread. The resulting Wigner volume chart is shown in Fig. 1(ii), and the relevant values are given in Table 1. The first FT rotates the chart by \( \pi/2 \) rad, see Fig. 1(iii), and the resultant signal passes through D2. The spatial frequency domain is again stretched (we note that this is actually the spatial domain of the input signal), and since the total energy again remains the same, the average energy density in phase space is further reduced; see Fig. 1(iv). The final FT rotates the chart once again as shown in Fig. 1(v). This 1D signal, with known extents in \( x \) and \( k_x \), is now captured. We see that it is important to match \( W_{IN} \) and \( W_{D1} \) and also \( \sqrt{\lambda f}B_{IN} + B_{D1} \) and \( W_{D1} \). Correct decryption of the signal follows a very similar pattern, in the exact reverse. Figures 1(ii)–1(v) in reverse order describe the evolution of a correctly decrypted signal. D2” now acts to despread the signal in the space domain, while D1” is not needed, since despread in the frequency domain does not affect the intensity of a decrypted signal.

Multiplexing is implemented by using a standard DRPE encryption system with different keys for each input image. Our demultiplexing system is shown in Fig. 2, made up of a standard DRPE decryption system except that we also use D1” to despread in \( k_x \). We also apply apertures in the decryption plane and its Fourier plane to remove all the energy of incorrectly decrypted signals except that part that lies inside the low-pass regions in the four dimensions. Figures 1(i)–1(v) in reverse describe the total evolution of demultiplexing. When decryption is performed by using an incorrect set of diffusers, the signal will not be de-
spread in \(x\) and \(k_x\) but rather will be even further spread in these domains, and the average energy density will be further reduced. In Fig. 1(iv) we show Wigner volume charts of a correctly and incorrectly decrypted signal. The extents of an incorrectly decrypted signal are given in Table 1.

A 2D encrypted signal will also have extents in \(y\) and \(k_y\), that may be determined by replacing the subscript \(x\) with \(y\) in Table 1. For 2D signals, the average energy density at any plane is given by

\[
E/\text{(Width } x\text{)}\text{(Width } y\text{)}/(\text{Width } k_x)/(\text{Width } k_y). \tag{5}
\]

We consider the case where we encrypt multiple images, each with different diffusers, and add the encrypted images together. To demultiplex we decrypt using the correct phase keys. The incorrectly decrypted signals will have been spread out in \(x\), \(y\), \(k_x\), and \(k_y\), and their average energy densities will be given by Eq. (5) and Table 1. The number of images that can be multiplexed together and recovered with and without loss is proportional to the phase space signal spreading. The total energy of each incorrectly decrypted signal present in the final output image Wigner volume is determined by multiplying the average energy density, \(\rho\), by the volume of interest, giving \(\rho B_{IN} W_{IN} B_{IN} W_{IN}\). Thus the average energy present due to incorrect decryption is given by \(\rho B_{IN} B_{IN}\). From this latter value the average error energy per bit can be determined.

Furthermore, each individual encryption is more robust than the traditional DPRE, since decryption requires knowledge of both the phase keys and filters used to despread and decrypt. If \(N\) signals are multiplexed together, to decode one we can expect an average error per bit of \(\text{AVE}_{IN}(N-1)\rho B_{IN} B_{IN}\) where \(\text{AVE}_{IN}\) represents the average energy of an input signal. Taking the latter value and the dynamic range of the input signals into account, one may determine the sufficient parameters for the diffusers (\(B_{D1}\) and \(B_{D1}\)) such that each image may be recovered with minimal error. We also note that a further limit is imposed on the number of images by the dynamic range of the recording medium. We believe that this limit may also be represented efficiently by using Wigner analysis, but that is beyond the scope of this Letter. In Fig. 3 we show the results of a simulation to multiplex and demultiplex three 128\(\times\)128 binary images. This case is clearly not optimal, since the input images have not been coded to evenly distribute the energy in the frequency domain. We chose \(B_{D1} = B_{D1}\), \(W_{IN} = \sqrt{IN} B_{D2}\), \(B_{IN} = B_{D1}\), \(W_{IN} = \sqrt{IN} B_{D2}\) so that the encrypted image had twice the width and bandwidth of the input, and therefore an incorrectly decrypted image had three times the width and bandwidth of the input. Thresholding was not applied to the results shown in Fig. 3. After thresholding the bit error rate for this case is \(<1\%\). The performance of the technique is currently being investigated for large numbers of inputs and different spreads. We acknowledge that a practical implementation will differ from the ideal simulation, and we are also pursuing experimental validation. Similar analyses may be applied to optimize the performance of architectures outlined in Refs. 6–14.

We acknowledge the support of the Irish Research Council for Science Engineering and Technology. B. Hennelly’s e-mail address is bryanh@cs.nuim.ie.

References